

Structural and optical characterization of transition metal dichalcogenide oolymorphs

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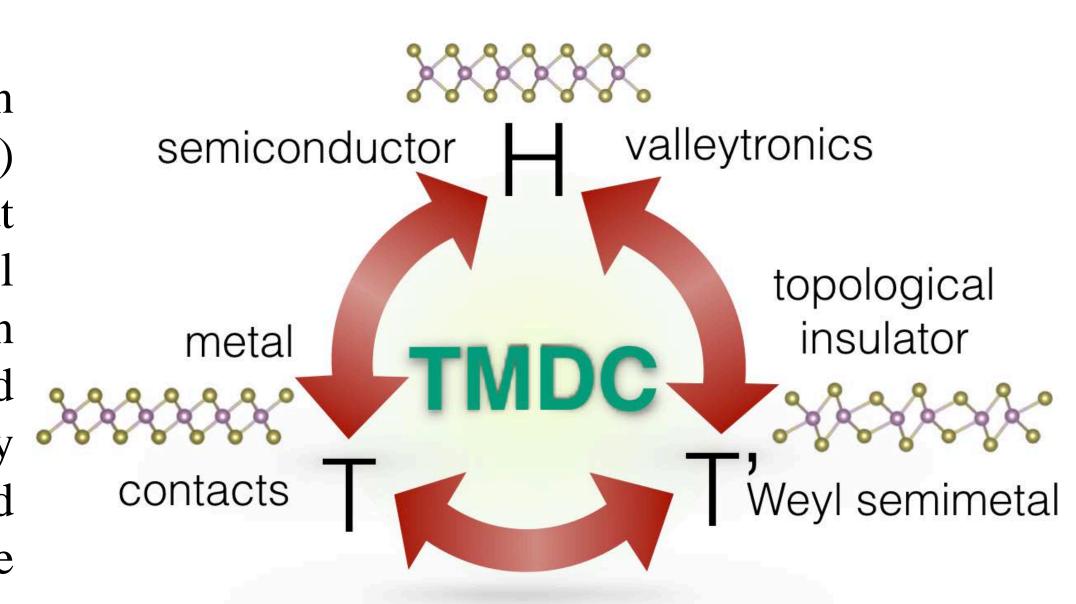
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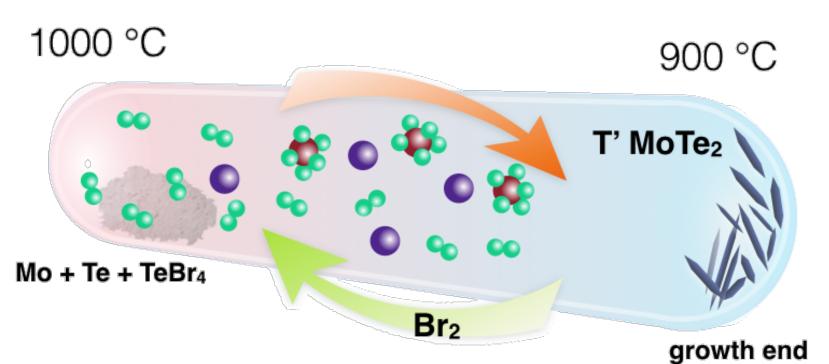
Introduction

Polymorphism of transition dichalcogenide (TMDC) recently attracted great attention due to its novel physical properties. Polarization and crystal-orientation resolved micro-Raman spectroscopy insightful and provides an noninvasive tool to study the TMDC polytypes.

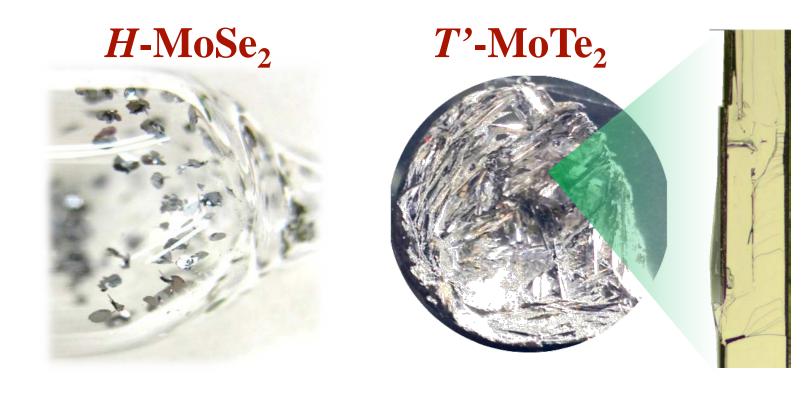


For H-TMDCs, we show that the scattered photons, even in monolayer H-TMDC, can have both helicity in `one' valley. For T'-TMDC, we report on the emergence of two low-wavenumber shear Raman bands in T'-MoTe₂ due to a phase transition from the monoclinic to the orthorhombic lattice, arising from inversion symmetry breaking.

Sample preparation

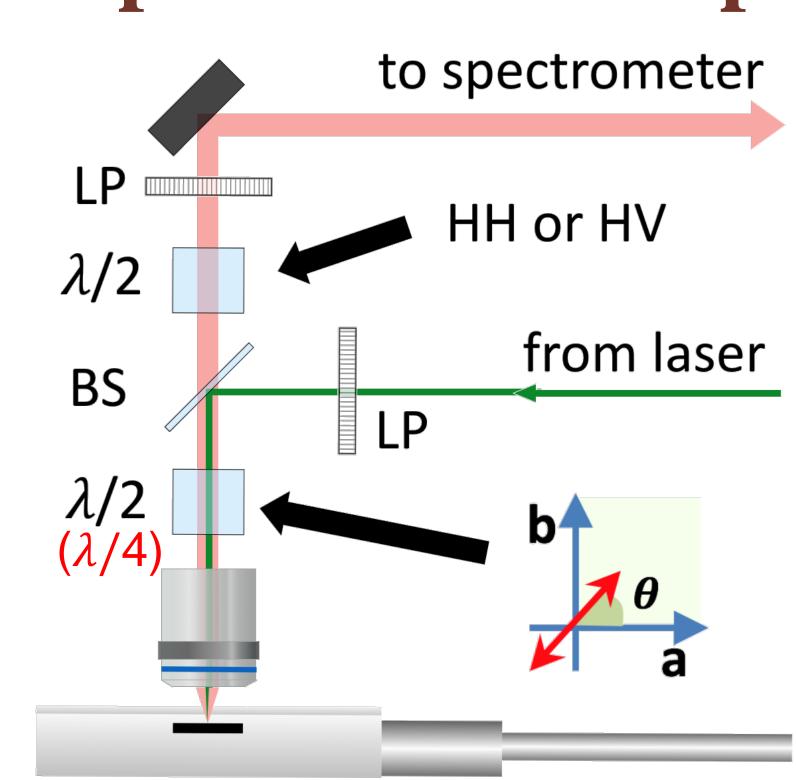


Schematic of CVT growth of T'-MoTe₂



Pictures of bulk *H*-MoSe₂ and *T*'-MoTe₂.

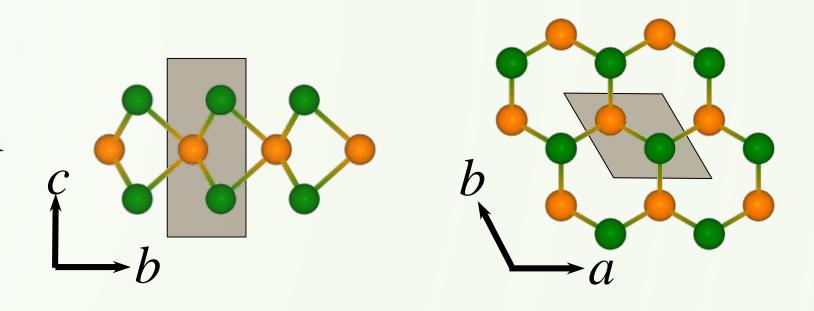
Experimental setup



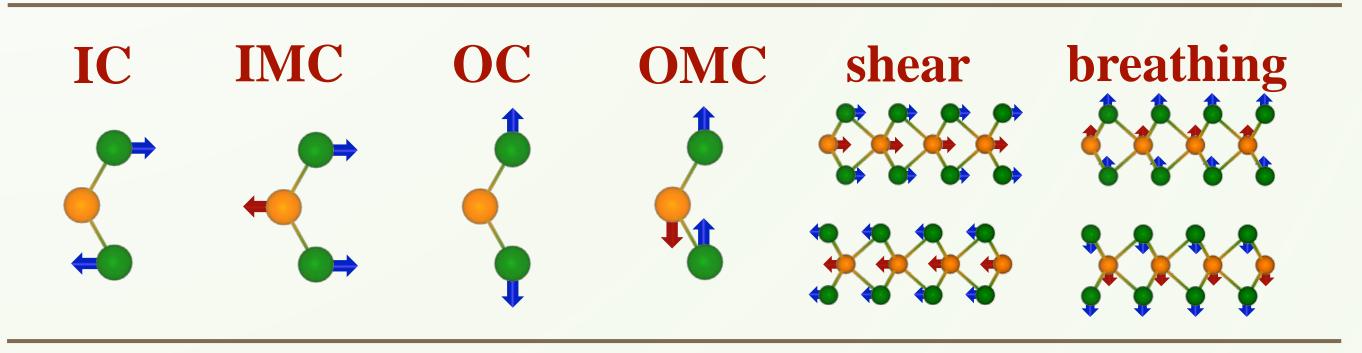
Experimental setup of polarization and crystal orientation resolved Raman spectroscopy.

H-TMDC

- M: transition metal
- X: chalcogen



H-TMDCs have a graphite-like structure with each graphene sheet replaced with an MX₂ triatomic layer. The zone center phonons can be classified into 6 types: 4 intralayer and 2 interlayer modes.

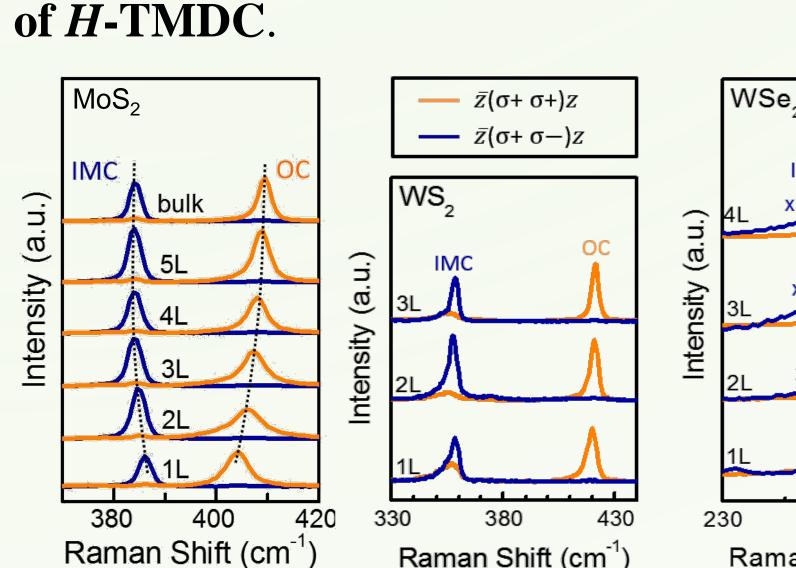


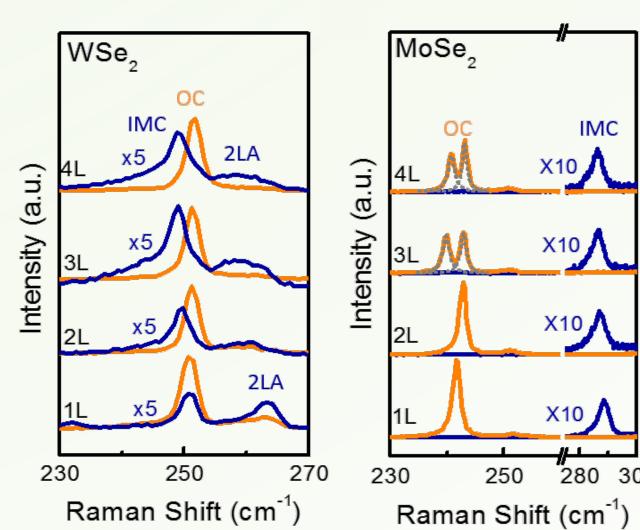
helicity PL

 $E_{\rm sc} = 2.54 \, {\rm eV}$

 $E_{\rm nh}$ (eV) 2.1 2.6

Distinct from helicity PL measurement which follows valley-photon selection rule, 🚊 1L helicity of Raman scattered light is sensitive to symmetry but **\\ \E** phonon insensitive to photon energy, a number of layer and types 1.7

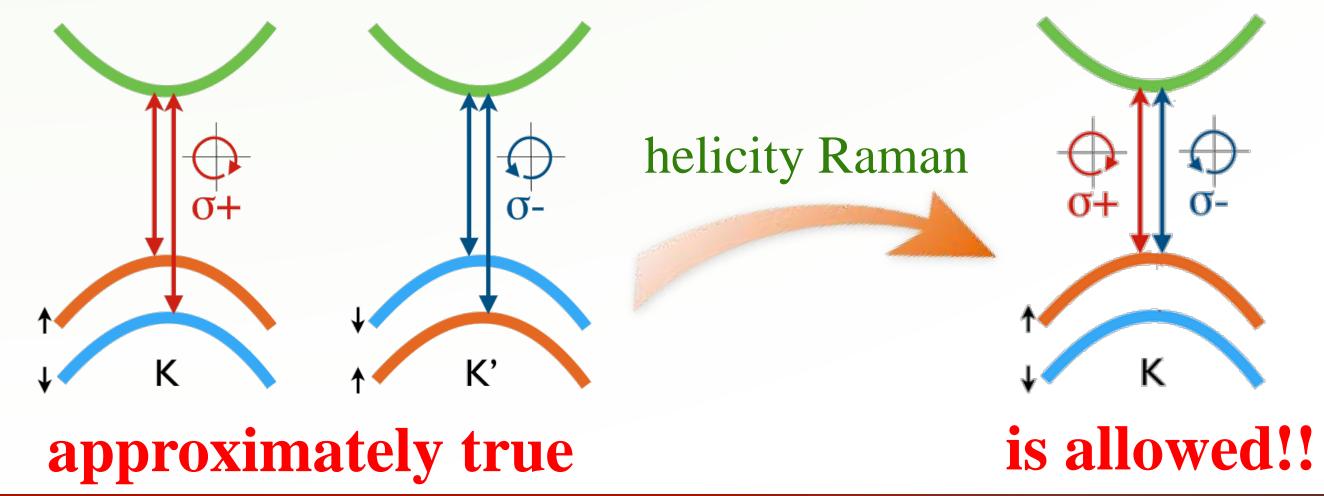




helicity Raman

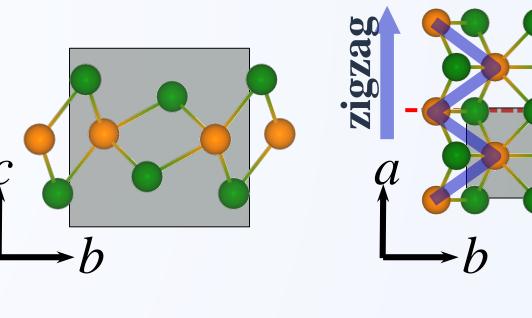
Raman Shift (cm⁻¹)

The selection rule of helicity of Raman is explained with Raman tensors calculation. Our results further show that the photon with opposite helicity can be emitted from the SAME valley, which indicates the valley-photon selection rule is approximately true.

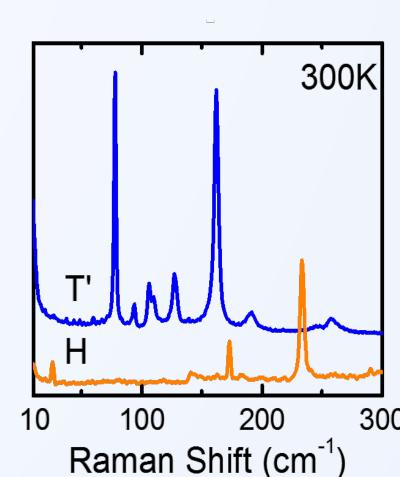


T'-TMDC

- M: transition metal
- X: chalcogen



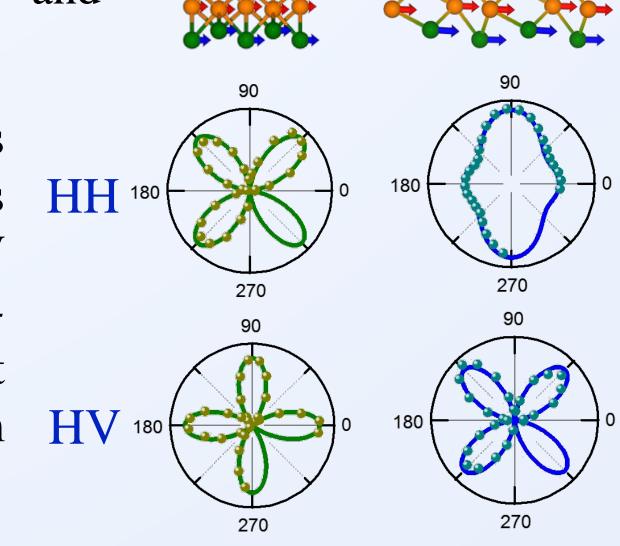
Due to strong interactions between transition metal atoms, T'-TMDCs form zigzag TM atomic chains (blue lines above) and exhibit distinct properties from H-TMDCs. The zone center phonons in T'-TMDCs can be categories: modes(z), vibrating along zigzag chains; and mirror-modes(m), vibrating in the mirror planes. (red dash line above)



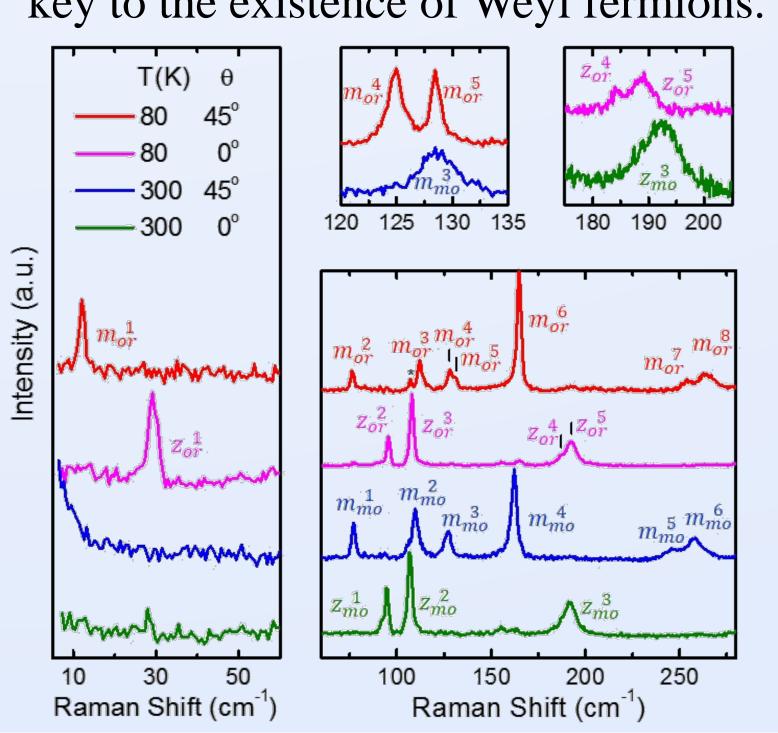
m-mode

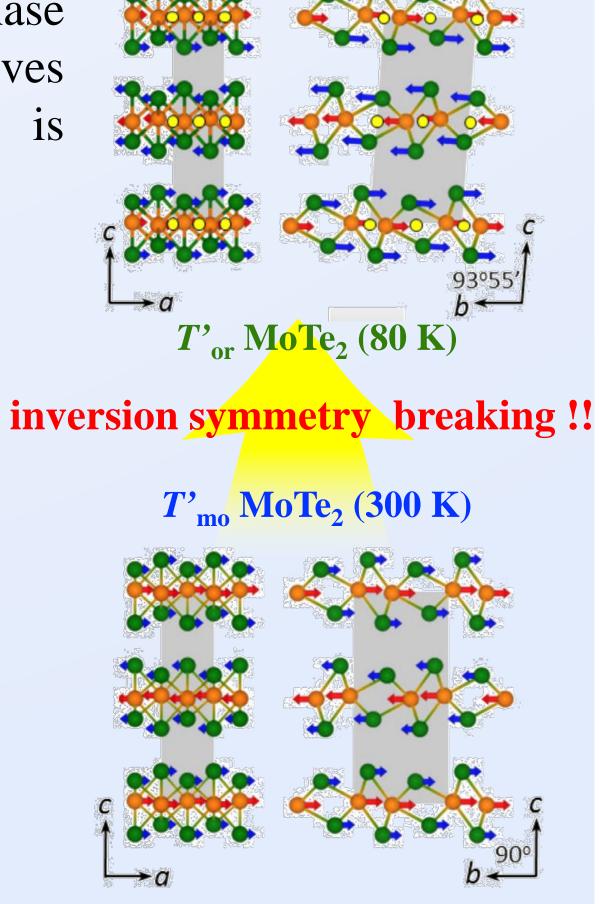
For bulk T'-MoTe₂, there are totally 36 zone-center phonons, including 24 mmode and 12 z-modes. Z-modes and mmodes have different symmetry, and show distinct angular dependence.

For HH configuration, z-modes show 4-fold symmetry; m-modes HH ¹⁸⁰ exhibit 2-fold symmetry. For HV configuration, both z-mode and mmodes display 4-fold symmetry but with their intensity maximum HV 180 shifted by 45 degree.



Four new Raman modes of T'-MoTe₂ are observed when cooled down. The phase transition from T'_{mo} to T'_{or} involves inversion-symmetry-breaking which is key to the existence of Weyl fermions.





Angle-dependent Raman cross section

The polarization and crystal orientation dependent Raman intensity can be calculated with the effective Raman tensor. The following is an example showing the calculation for m-mode and z-mode in T'-MoTe₂.

$$I = A| < \epsilon_i | R^T \cdot \Re \cdot R | \epsilon_o > |^2$$

$$I_{HH}^{z}(\theta) = |a\cos \theta + e\sin \theta|$$

$$I_{HH}^{z}(\theta) = |g|^2 \sin^2(2\theta)$$

$$R(A_1) = \begin{pmatrix} d & 0 & 0 \\ 0 & e & 0 \\ 0 & 0 & f \end{pmatrix}, R(A_2) = \begin{pmatrix} 0 & g & 0 \\ g & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

$$I_{HV}^{z}(\theta) = \left| \frac{d - e}{2} \right|^2 \sin^2(2\theta),$$

$$I_{HV}^{z}(\theta) = |a|^2 \cos^2(2\theta),$$

$$I_{HV}^{z}(\theta) = |a|^2 \cos^2(2\theta),$$

$$I_{HH}^m(\theta) = \left| d\cos^2 \theta + e\sin^2 \theta \right|^2,$$

$$I_{HH}^z(\theta) = \left| g \right|^2 \sin^2(2\theta).$$

$$I_{HV}^{z}(\theta) = |g|^{2} \cos^{2}(2\theta).$$